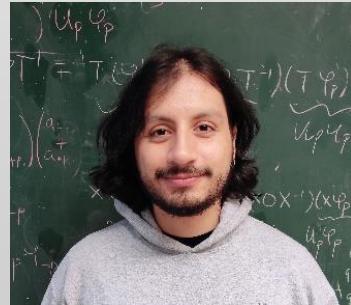


# Fragility of spectral flow in topological insulators

**Matthew S. Foster**

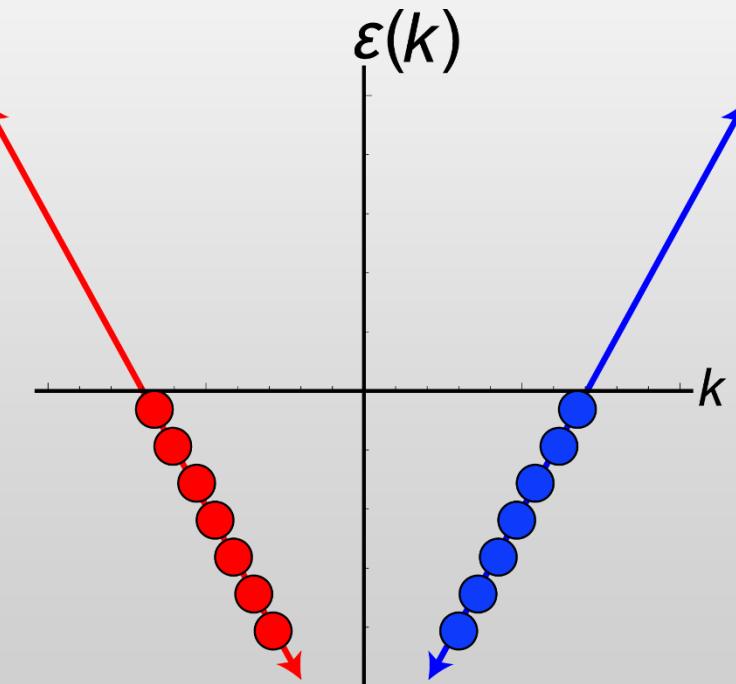
**December 12, 2023**



A. Altland, P. Brouwer, J. Dieplinger, M. S. Foster, M. Moreno-Gonzalez, and L. Trifunovic  
[arXiv:2308.12931](https://arxiv.org/abs/2308.12931)

## Chiral anomaly and spectral flow

- **Massless Dirac Hamiltonian in  $d = 1$ :**  $\hat{h} = -iv_F \hat{\sigma}^3 \frac{d}{dx}$
- **This is a 2-band semimetal**



## Chiral anomaly and spectral flow

- **Massless Dirac Hamiltonian in  $d=1$ :**  $\hat{h} = -iv_F \hat{\sigma}^3 \frac{d}{dx}$
- **Couple 1+1-D Dirac to electromagnetism: Axial anomaly**
  1. Electric “2-current” is conserved  $\partial_t \rho + \partial_x J = 0$
  2. Axial (swap) 2-current is not

$$\partial_t \left( \frac{J}{v_F} \right) + \partial_x (v_F \rho) = \frac{2e^2}{h} E$$

Schwinger 1962

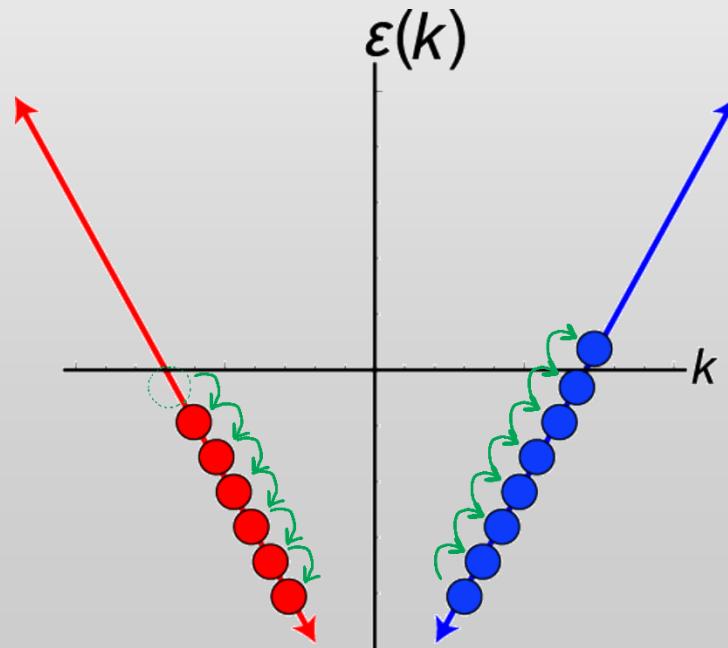
## Chiral anomaly and spectral flow

- Axial current non-conservation: Material realizations

$$\frac{1}{v_F} \frac{dJ}{dt} = e \frac{d}{dt} (n_R - n_L) = \frac{2e^2}{h} E$$

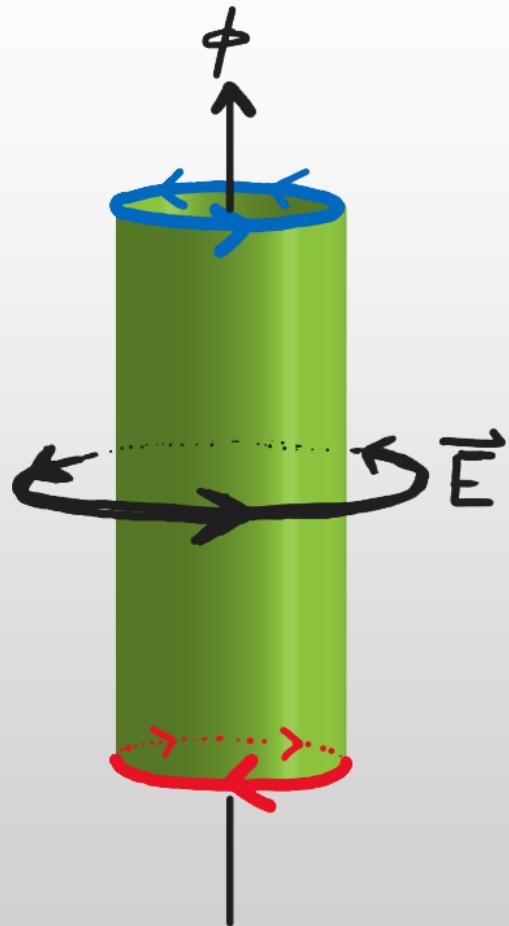
1. Newton's 2<sup>nd</sup> law, 1D wire:

**Left-movers convert to right-movers (through the bottom of the band, not shown)**



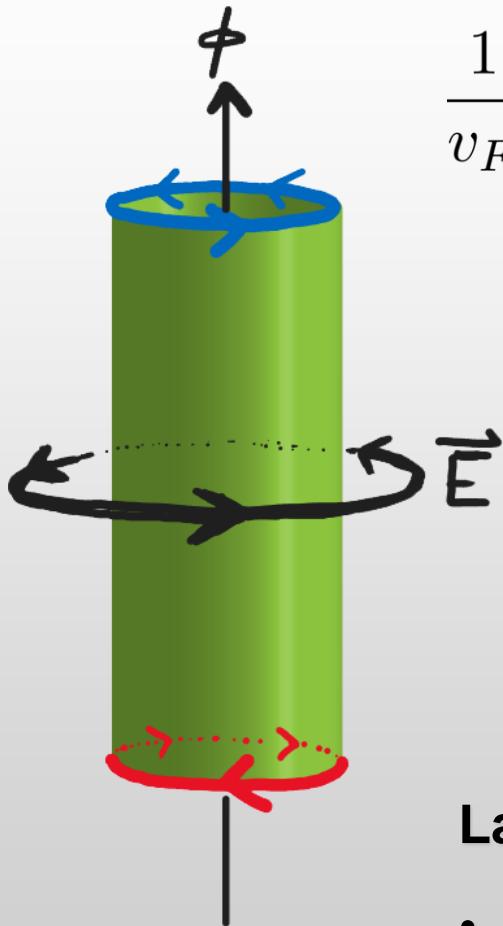
## Chiral anomaly and spectral flow

What if we could physically separate right- and left-movers?



## Chiral anomaly and spectral flow

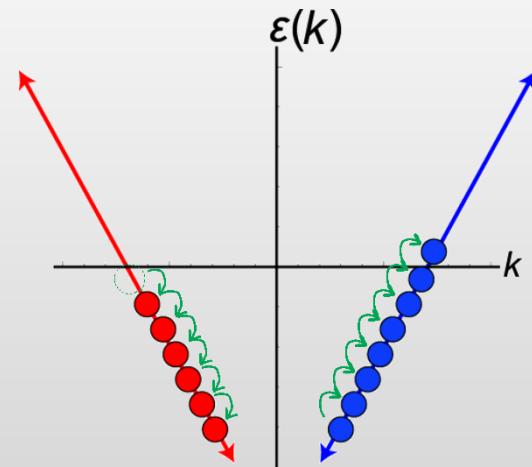
What if we could physically separate right- and left-movers?



$$\frac{1}{v_F} \frac{dJ}{dt} = e \frac{d}{dt} (n_R - n_L) = \frac{2e^2}{h} E$$

$$\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$$

$$\Delta N_R - \Delta N_L = \frac{2e}{hc} \Delta \phi$$



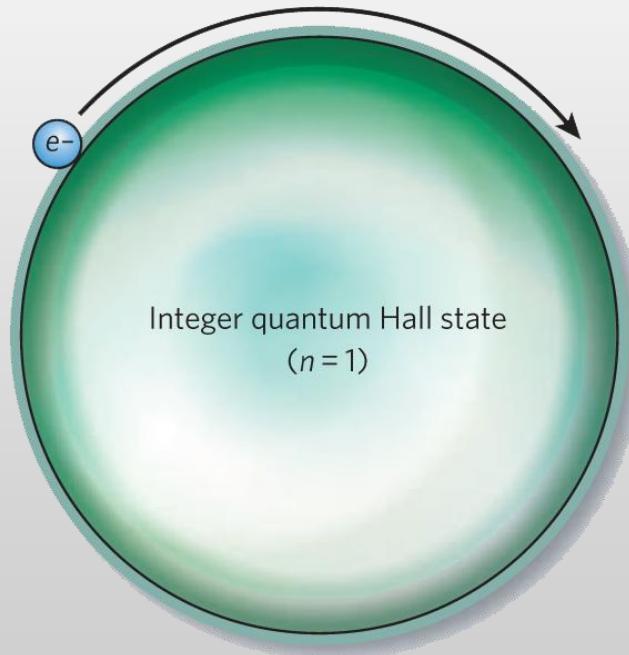
Laughlin 1981, Halperin 1982:

- Pump from right- to left- edges:
- “**Spectral Flow Principle**” for topological matter!

# Anderson Localization is ubiquitous in 1,2-D...Topological exceptions

## 1. 1D chiral edge states of a quantum Hall droplet

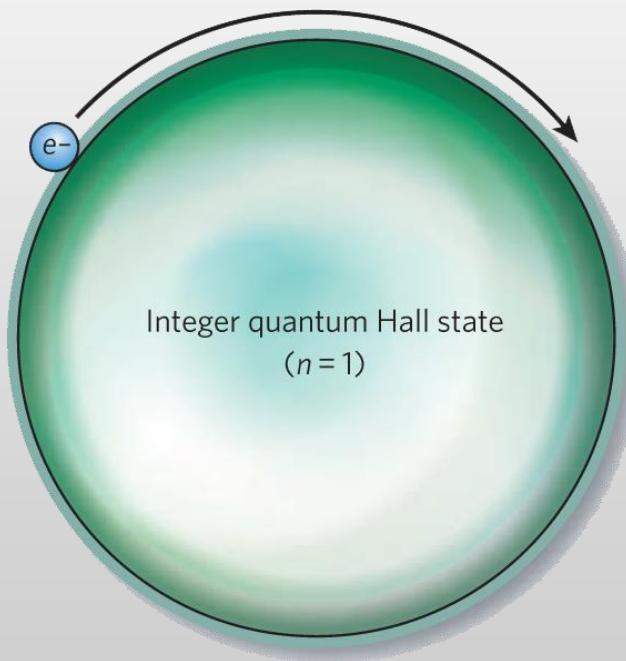
$$\hat{h} = -iv_F \frac{d}{dx}$$



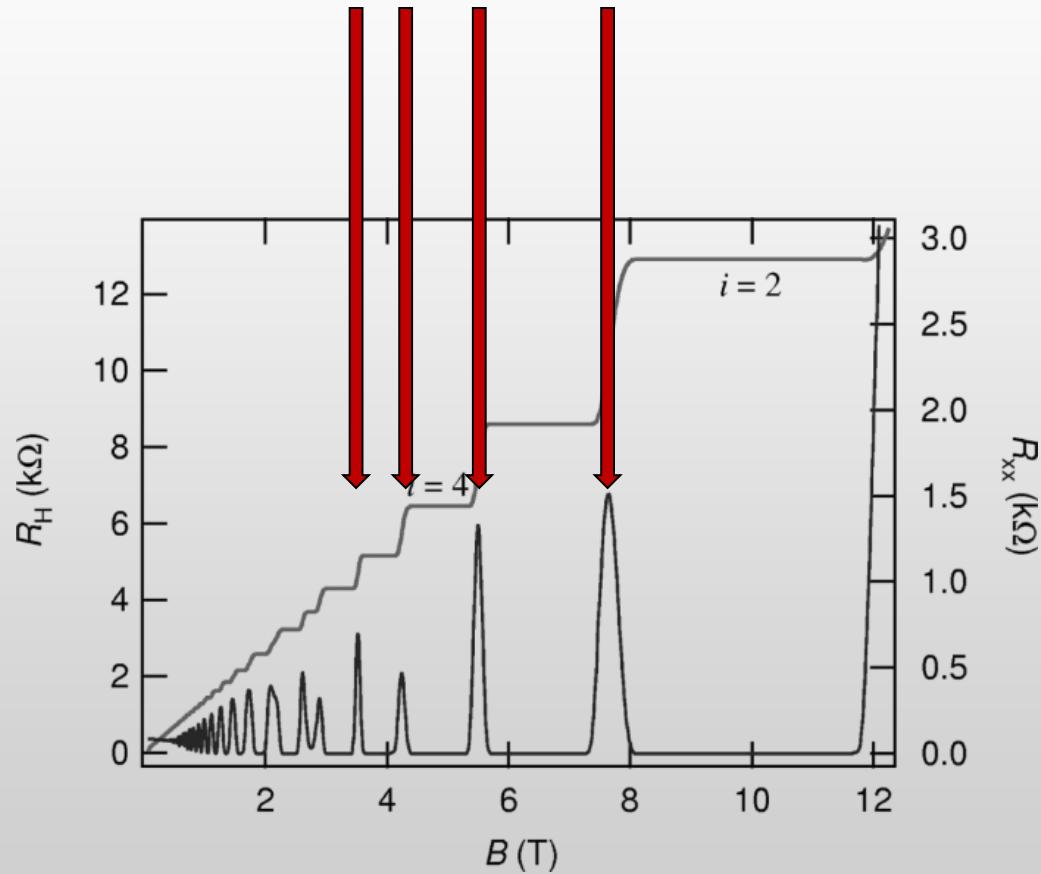
# Anderson Localization is ubiquitous in 1,2-D...Topological exceptions

1. 1D chiral edge states of a quantum Hall droplet

$$\hat{h} = -iv_F \frac{d}{dx}$$



2. At the quantum phase transition between Hall plateaux

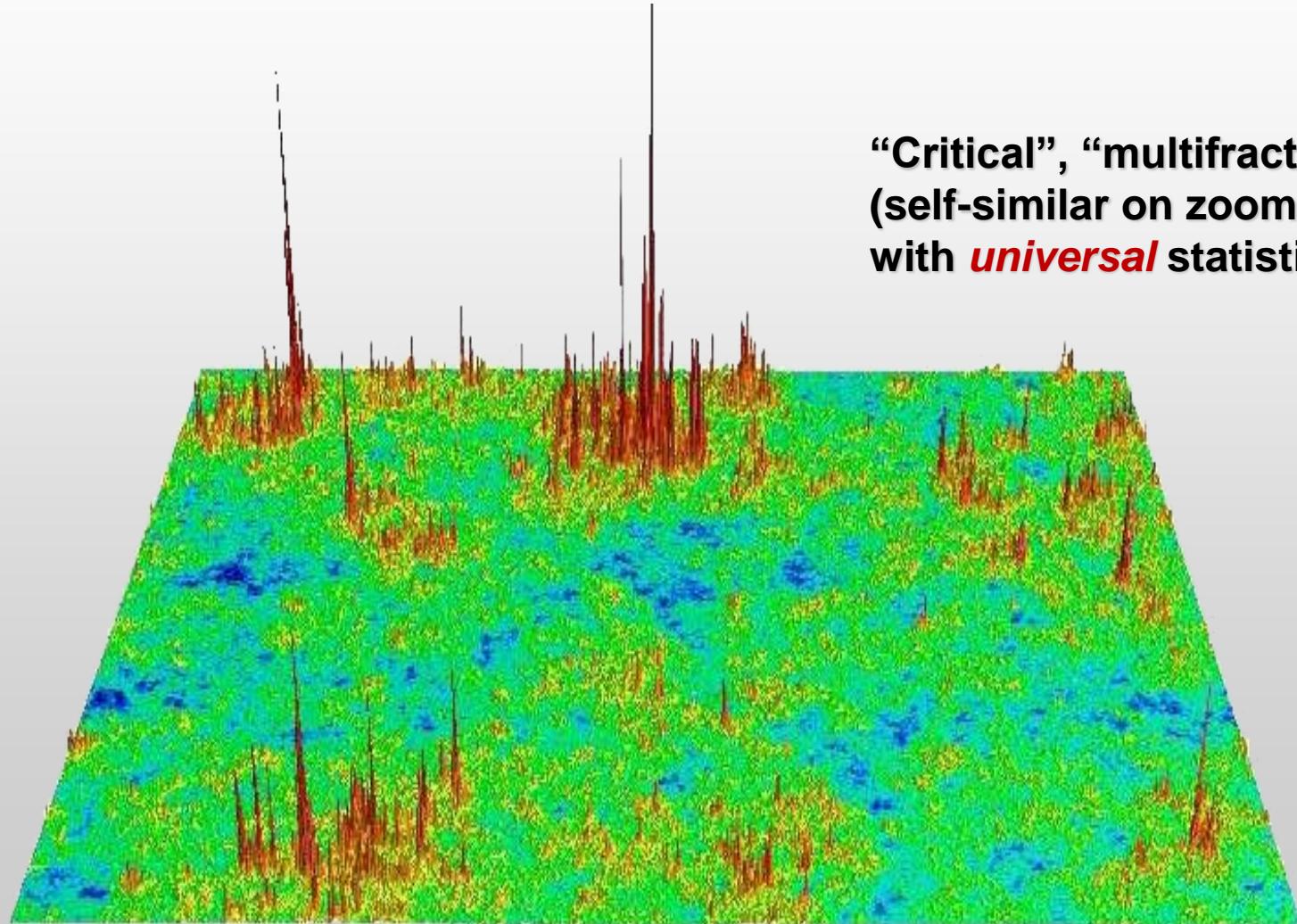


## **Anderson Localization is ubiquitous in 1,2-D...Topological exceptions**

**What does a non-localized, plateau-transition wave function look like?**

# Anderson Localization is ubiquitous in 1,2-D...Topological exceptions

What does a non-localized, plateau-transition wave function look like?



F. Evers

## 2D Dirac fractality: ...Topology?

- Where else can Anderson localization be avoided in 2D?
- 2+1-D Dirac fermions with **gauge** disorder...but no “natural” realization in graphene, TI surface, etc.

$$\hat{h} = v_F \left\{ \hat{\sigma}^1 [-i\partial_x + A_x(\mathbf{r})] + \hat{\sigma}^2 [-i\partial_y + A_y(\mathbf{r})] \right\}$$

## 2D Dirac fractality: ...Topology?

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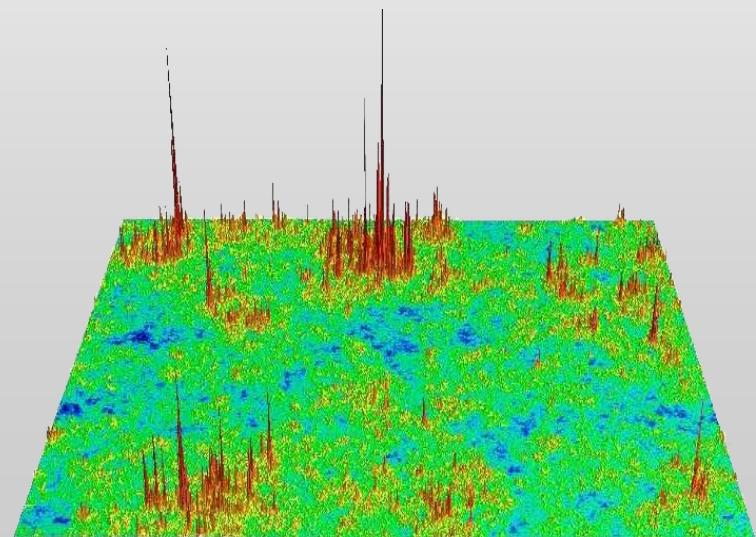
- “Quenched” QED or QCD: disorder can be abelian or non-abelian (for multiple cones)
- Exact solution (conformal field theory):  
**Zero-energy eigenstate is always extended (multifractal)**

Ludwig, Fisher, Shankar, Grinstein (1994)

Nersesyan, Tsvelik, Wenger (1995)

Mudry, Chamon, Wen (1996)

Caux, Kogan, Tsvelik (1996)



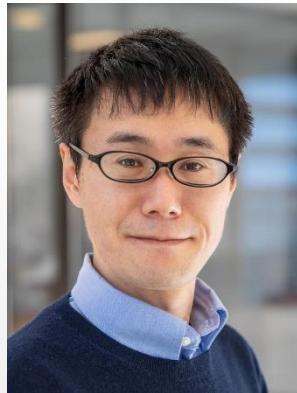
# 3D Topological superconductors: Discovered by Dirac fractality!



From the 2018 film “Annihilation”



# 3D Topological superconductors: Discovered by Dirac fractality!



$$\hat{h} = v_F \left\{ \hat{\sigma}^1 [-i\partial_x + A_x(\mathbf{r})] + \hat{\sigma}^2 [-i\partial_y + A_y(\mathbf{r})] \right\}$$

**Schnyder, Ryu, Furusaki, Ludwig (2008):**

- These are **surface states of new topological phases in 3D!**
- Bulk of these new phases: **Topological superconductors!**

# 3D Topological superconductors: Discovered by Dirac fractality!

Class	$T$	$P$	$S$	Spin sym.	$d = 2$	$d = 3$	Topological realization	Replicated fermion $N\sigma M$
C	0	-1	0	SU(2)	$2\mathbb{Z}$	...	SQHE (2D $d + id$ TSC)	$\text{Sp}(4n)/\text{U}(2n)$
A (unitary)	0	0	0	U(1)	$\mathbb{Z}$	...	IQHE	$\text{U}(2n)/\text{U}(n) \otimes \text{U}(n)$
D	0	+1	0	...	$\mathbb{Z}$	...	TQHE (2D $p + ip$ TSC)	$\text{O}(2n)/\text{U}(n)$
CI	+1	-1	1	SU(2)	...	$2\mathbb{Z}$	3D TSC	$\text{Sp}(4n) \otimes \text{Sp}(4n)/\text{Sp}(4n)$
AIII	0	0	1	U(1)	...	$\mathbb{Z}$	3D TSC, chiral TI	$\text{U}(2n) \otimes \text{U}(2n)/\text{U}(2n)$
DIII	-1	+1	1	...	$\mathbb{Z}_2$	$\mathbb{Z}$	3D TSC ( ${}^3\text{He}-B$ )	$\text{O}(2n) \otimes \text{O}(2n)/\text{O}(2n)$
AI (orthogonal)	+1	0	0	SU(2)	...	...	...	$\text{Sp}(4n)/\text{Sp}(2n) \otimes \text{Sp}(2n)$
AII (symplectic)	-1	0	0	...	$\mathbb{Z}_2$	$\mathbb{Z}_2$	2D, 3D TIs	$\text{O}(2n)/\text{O}(n) \otimes \text{O}(n)$
BDI	+1	+1	1	SU(2)	...	...	...	$\text{U}(2n)/\text{Sp}(2n)$
CII	-1	-1	1	...	...	$\mathbb{Z}_2$	3D chiral TI	$\text{U}(2n)/\text{O}(2n)$

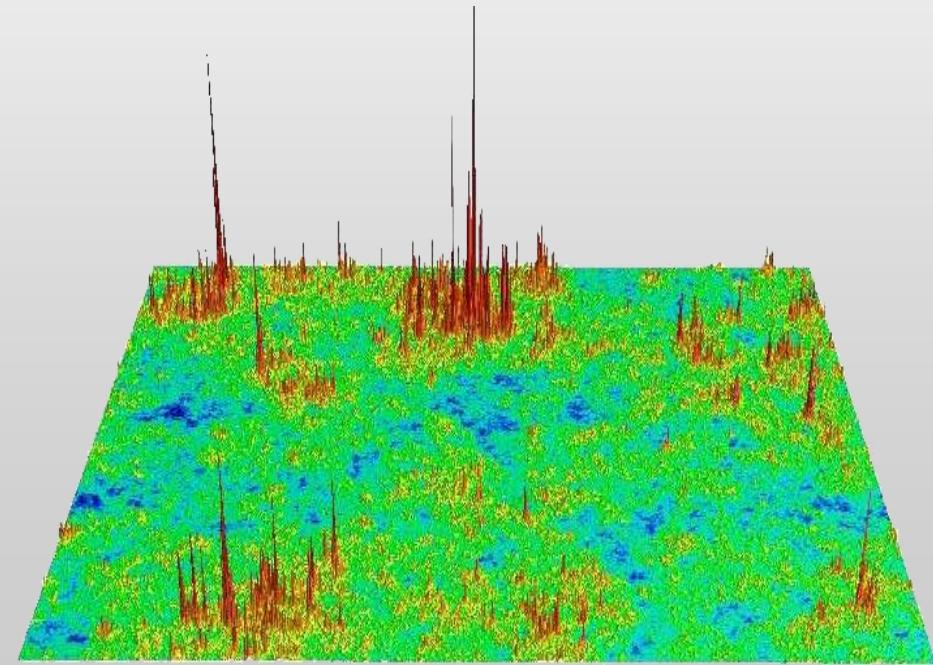
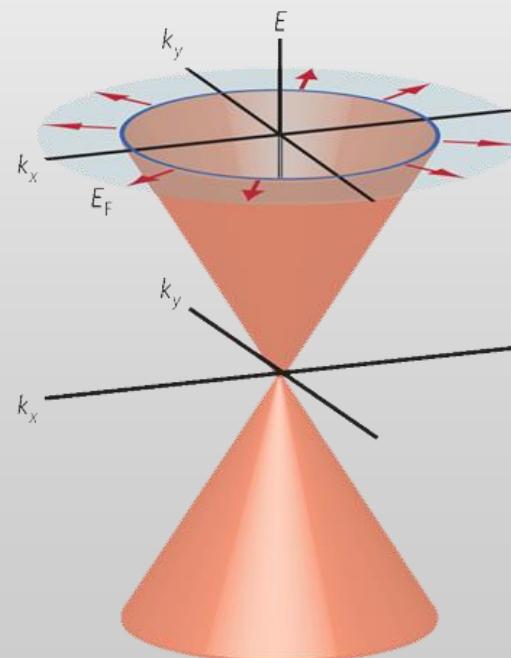
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# 3D Topological superconductors: Discovered by Dirac fractality!

- Zero-energy state is “topologically protected” (not localized)
- ...what about all of the others?

$$\hat{h} = v_F \left( -i\hat{\sigma}^1 [-i\partial_x + A_x(\mathbf{r})] - i\hat{\sigma}^2 [-i\partial_y + A_y(\mathbf{r})] \right)$$



Class	$T$	$P$	$S$	Spin sym.	$d = 2$	$d = 3$	Topological realization	Replicated fermion NL $\sigma$ M
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BDI	+1	+1	1	SU(2)	...	...	...	$\text{U}(2n)/\text{Sp}(2n)$
CII	-1	-1	1	...	...	$\mathbb{Z}_2$	3D chiral TI	$\text{U}(2n)/\text{O}(2n)$

# 3D topological superconductors: Spectrum-wide quantum criticality

Class	$T$	$P$	$S$	Spin sym.	$d = 2$	$d = 3$	Topological realization	Replicated fermion NL $\sigma$ M
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AI (orthogonal)	+1	0	0	SU(2)	...	...	...	$Sp(4n)/Sp(2n) \otimes Sp(2n)$
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BDI	+1	+1	1	SU(2)	...	...	...	$U(2n)/Sp(2n)$
CII	-1	-1	1	...	...	$\mathbb{Z}_2$	3D chiral TI	$U(2n)/O(2n)$

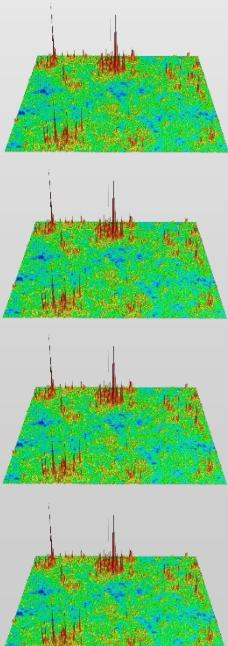
## Spectrum-wide quantum criticality conjecture:

Generic surface states of 3D topological superconductors are critical, **statistically identical to plateau transition states** in 2D quantum Hall effects

Ghorashi,  
Liao, Foster  
PRL (2018)

# 3D topological superconductors: Spectrum-wide quantum criticality

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BDI	+1	+1	1	SU(2)	...	...	...	$U(2n)/Sp(2n)$
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**Spectrum-wide quantum criticality conjecture:**

Generic surface states of 3D topological superconductors are critical, statistically identical to plateau transition states in 2D quantum Hall effects



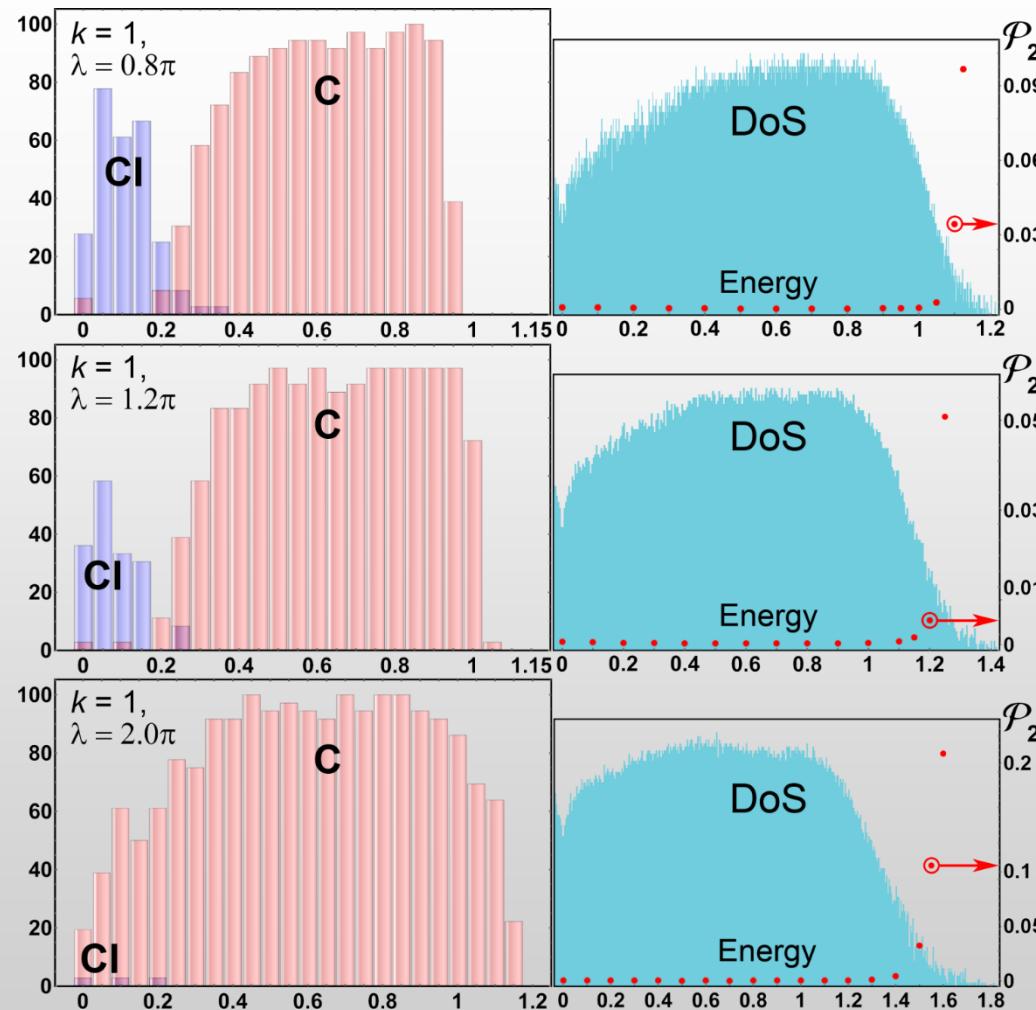
# Numerical study: 2D Dirac surface states of 3D TSCs



1. Ghorashi, Liao, Foster PRL (2018)
2. Sbierski, Karcher, Foster PRX (2020)
3. Ghorashi, Karcher, Davis, Foster PRB (2020)

**Review:**  
Karcher and Foster  
Ann Phys (2021)

# Numerical study: 2D Dirac surface states are protected at ALL energies!



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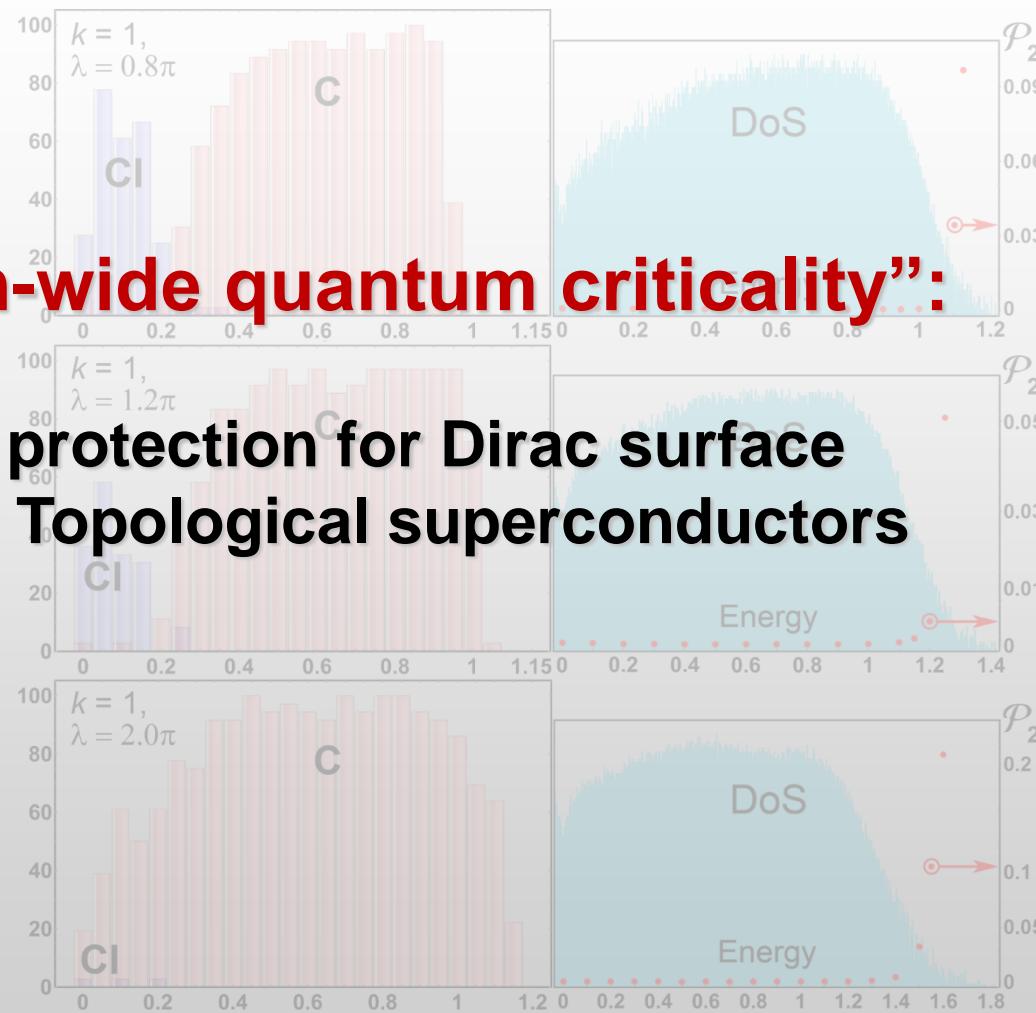
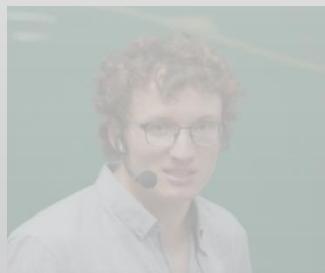
# Numerical study: 2D Dirac surface states are protected at ALL energies!



**“Spectrum-wide quantum criticality”:**

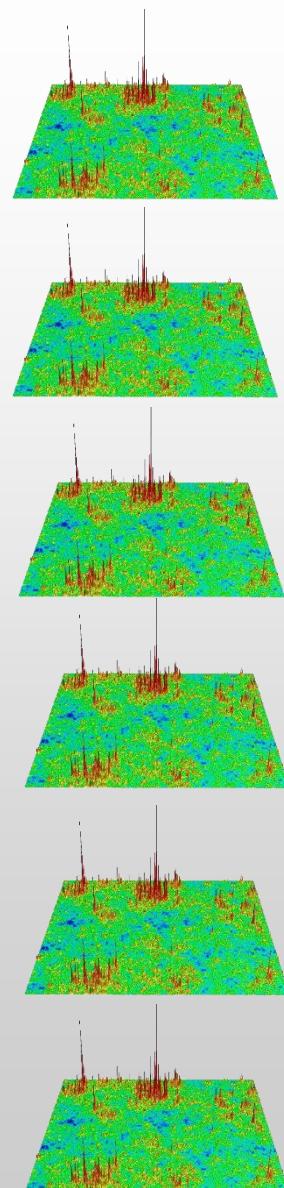


**Topological protection for Dirac surface states of 3D Topological superconductors**



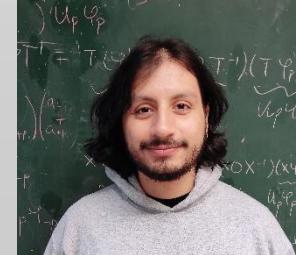
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## ... Enter the “UFO” (Fragmenting potential)

- We consider a certain lattice model of an AIII topological phase
- There is a **strange perturbation** that we can add uniformly to the surface (possessing Dirac surface states)

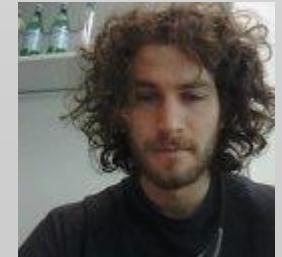
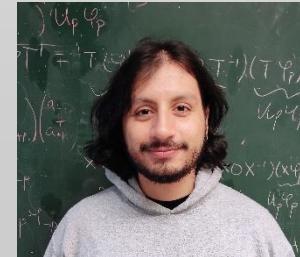


A. Altland, P. Brouwer, J. Dieplinger, M. S. Foster, M. Moreno-Gonzalez, and L. Trifunovic  
**arXiv:2308.12931**

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  1. It preserves the defining symmetry of the class

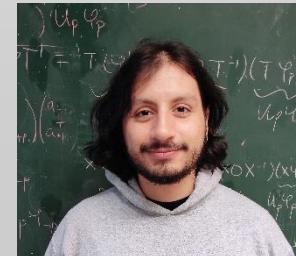
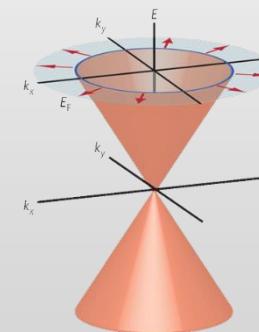
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## ... Enter the “UFO” (Fragmenting potential)

- We consider a certain lattice model of an AIII topological phase
- There is a **strange perturbation** that we can add uniformly to the surface (possessing Dirac surface states)
  1. It preserves the defining symmetry of the class
  2. It projects to zero in the Dirac description of the surface

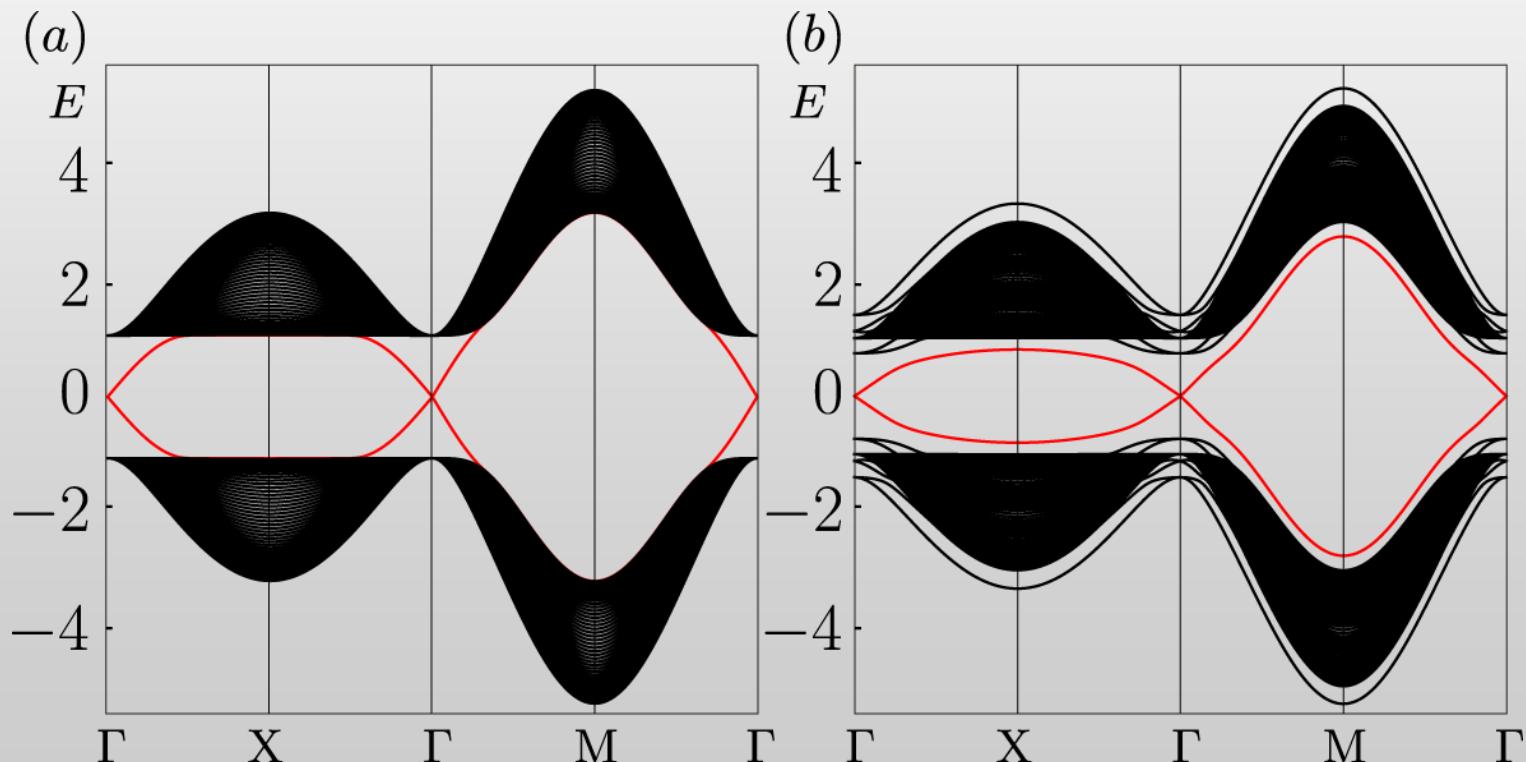
$$\hat{h} = v_F \{ \hat{\sigma} \cdot [-i\nabla + \mathbf{A}(\mathbf{r})] \}$$



## ... Enter the “UFO” (Fragmenting potential)

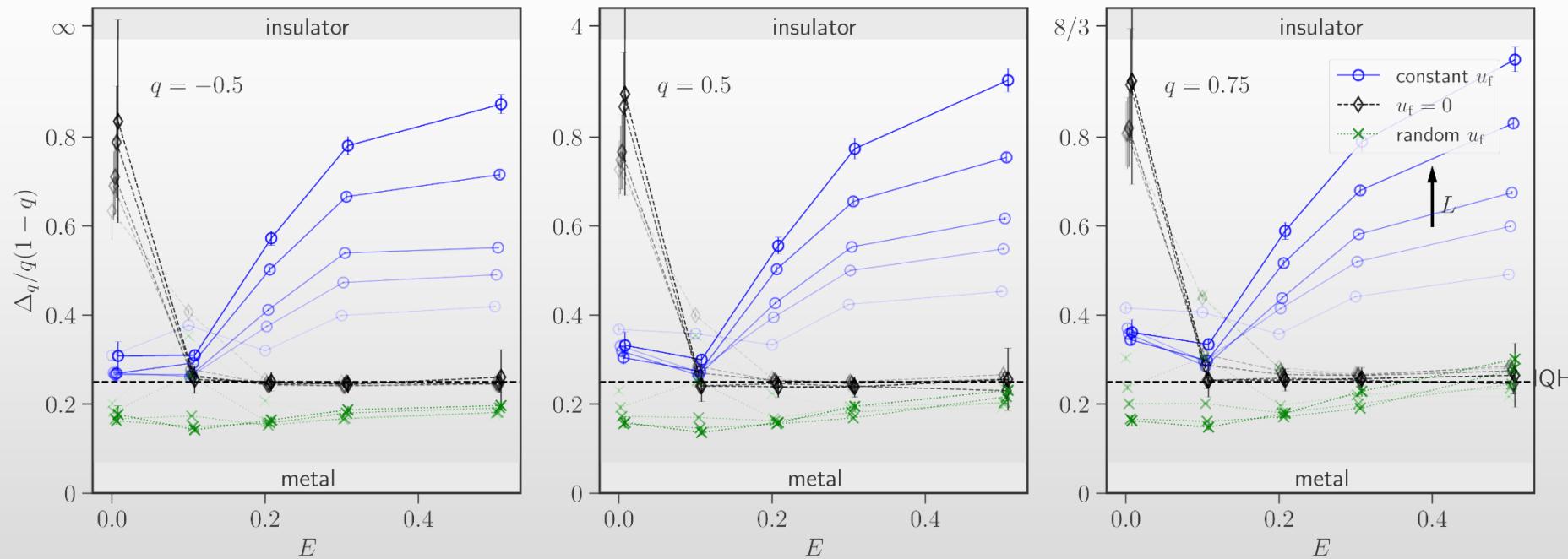
- We consider a certain lattice model of an AIII topological phase
- There is a **strange perturbation** that we can add uniformly to the surface (possessing Dirac surface states)

### 3. ...and, it **interrupts spectral flow!**



## ... Enter the “UFO” (Fragmenting potential)

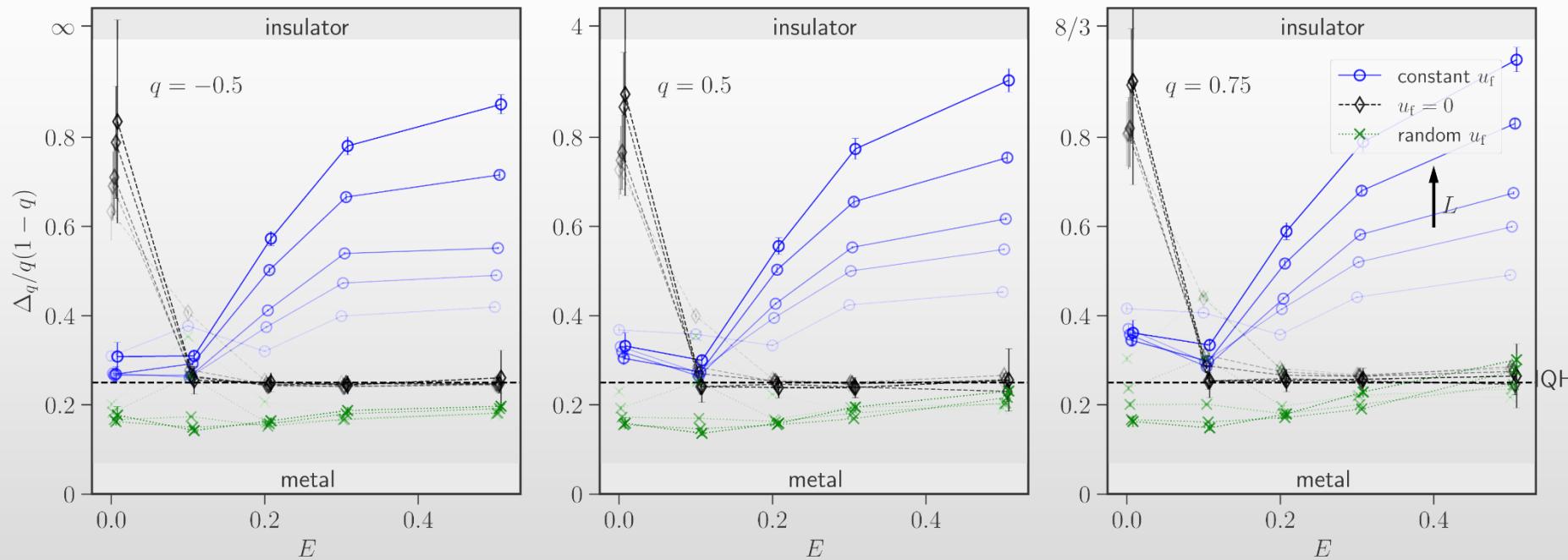
### 4. ...and it Anderson localizes almost all surface states!



- **Black data: No UFO. Spectrum-wide criticality (no localization)**
- **Blue data: Uniform UFO. Localization except at  $E = 0$ .**
- **Green data: Random UFO...spectrum-wide critical again!**

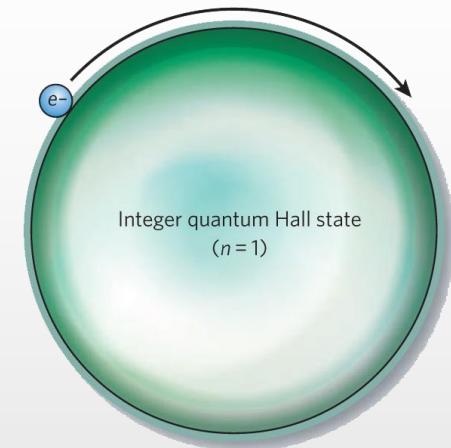
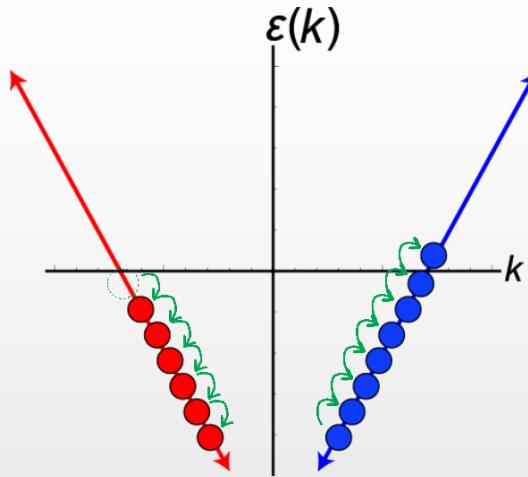
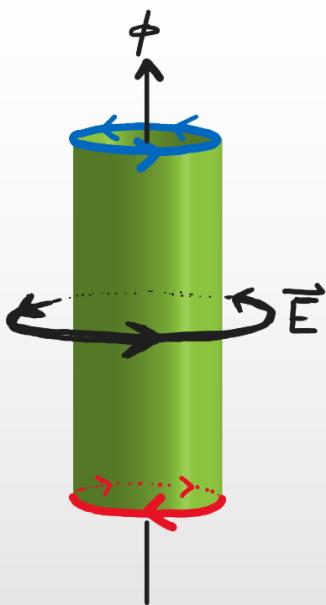
## ... Enter the “UFO” (Fragmenting potential)

### 4. ...and it Anderson localizes almost all surface states!



- All surface states conduct without the UFO
- Almost none conduct with it...unless *random* with zero average
- Dirac equation can't tell the difference (UFO projects to zero!)

# Why doesn't Dirac work? Missing quantum geometry



Back to Quantum Hall...

- Edge view: Axial anomaly of 1+1-D Dirac equation
- Bulk view: Topological winding number  $W$  due to **Berry curvature**

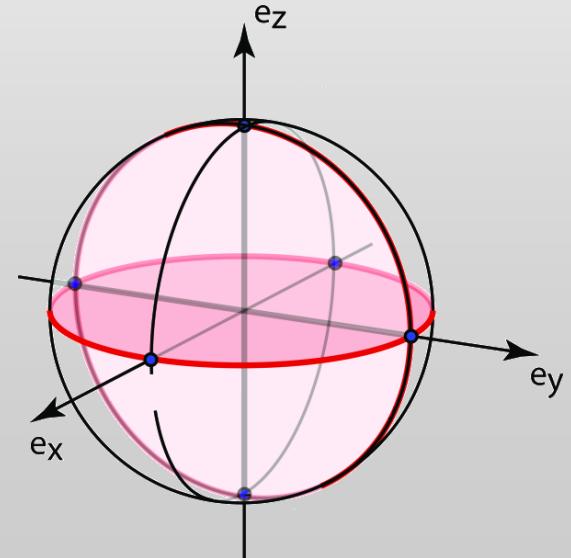
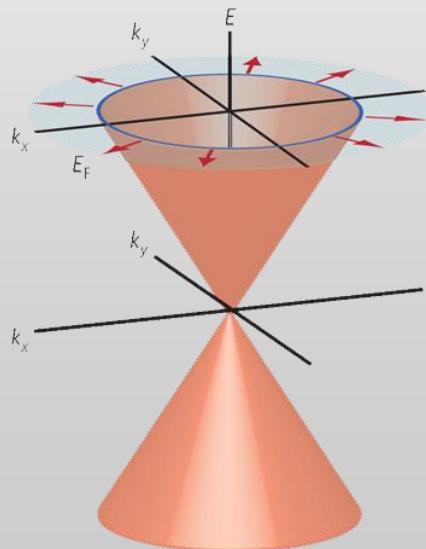
$$\Omega_{\mathbf{k}} = i \langle d\alpha_{\mathbf{k}} | \wedge d\alpha_{\mathbf{k}} \rangle, \quad W = \frac{1}{2\pi} \int_{\mathbf{k}} \Omega_{\mathbf{k}}$$

# Why doesn't Dirac work? Missing quantum geometry

## The “UFO” (fragmenting potential)

- Induces *surface* Berry curvature “in the sky,”  
*Not captured by the Dirac equation for the surface (Berry-flat)*

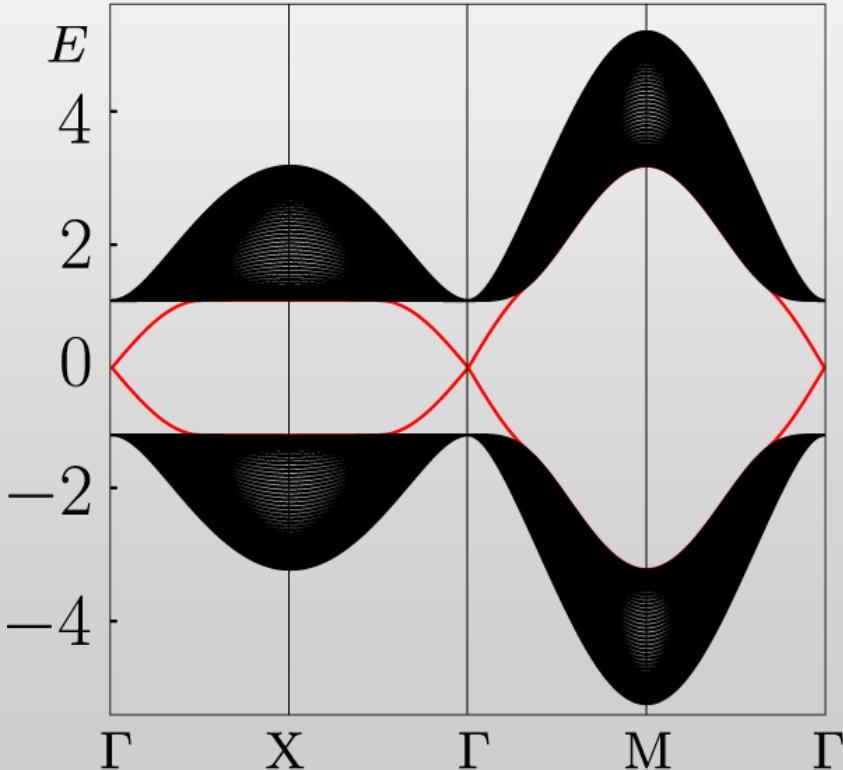
$$\psi_{\text{Surf}}^{(\text{Dirac})}(k_x, k_y) = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ \pm \frac{(k_x + ik_y)}{|\mathbf{k}|} \end{bmatrix}$$



# Why doesn't Dirac work? Missing quantum geometry

## The “UFO” (fragmenting potential)

- Induces *surface Berry curvature “in the sky,”*  
*Not captured by the Dirac equation for the surface (Berry-flat)*



$$\psi_{\text{Surf}}^{(\text{Lattice})}(k_x, k_y) = \begin{bmatrix} \psi_1(\mathbf{k}) \\ \psi_2(\mathbf{k}) \\ \psi_3(\mathbf{k}) \\ \psi_4(\mathbf{k}) \end{bmatrix} \rightarrow \begin{bmatrix} 1 \\ \pm \frac{(k_x + ik_y)}{|\mathbf{k}|} \\ \pm 1 \\ \frac{(k_x + ik_y)}{|\mathbf{k}|} \end{bmatrix}$$

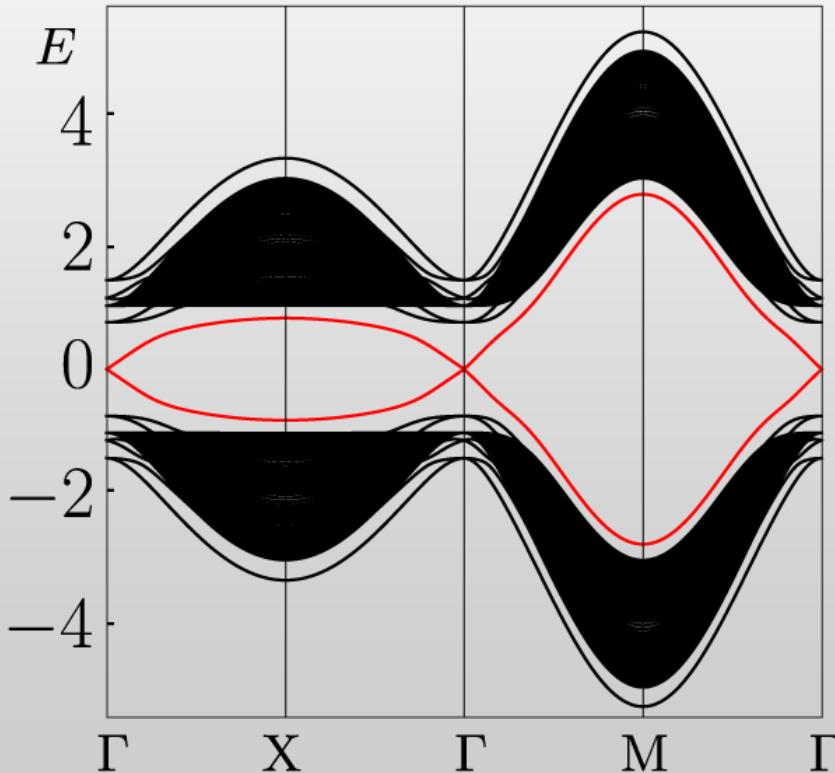
No UFO:

4 components deviate from Dirac,  
but remain Berry-flat

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- Induces *surface Berry curvature “in the sky,”*  
*Not captured by the Dirac equation for the surface (Berry-flat)*



$$\psi_{\text{Surf}}^{(\text{Lattice})}(k_x, k_y) = \begin{bmatrix} \psi_1(\mathbf{k}) \\ \psi_2(\mathbf{k}) \\ \psi_3(\mathbf{k}) \\ \psi_4(\mathbf{k}) \end{bmatrix} \rightarrow \begin{bmatrix} 1 \\ \pm \frac{(k_x + ik_y)}{|\mathbf{k}|} \\ \pm 1 \\ \frac{(k_x + ik_y)}{|\mathbf{k}|} \end{bmatrix}$$

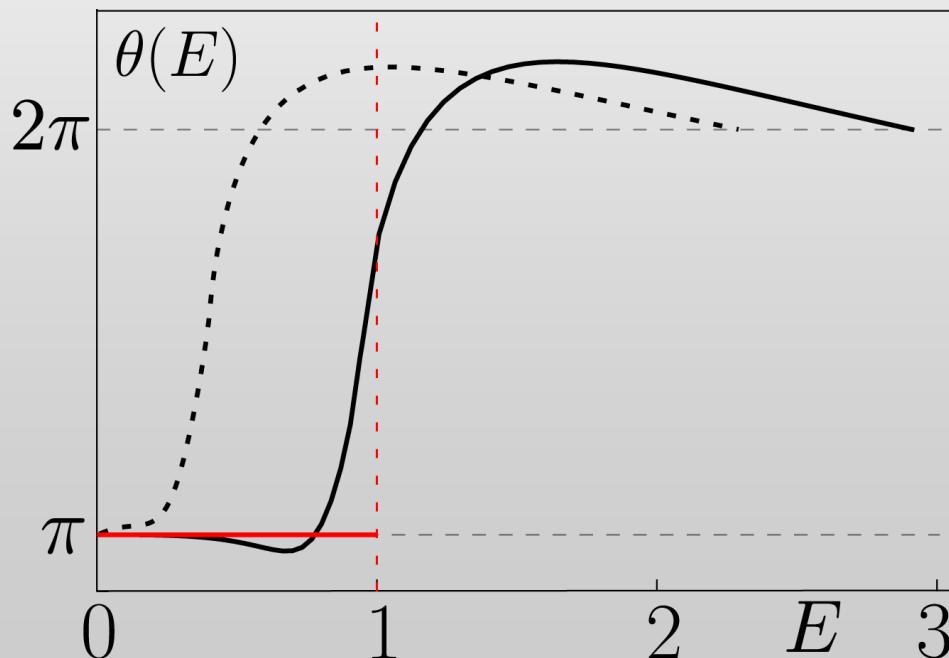
With UFO:

4 components deviate from Dirac,  
and develop Berry curvature!

# Why doesn't Dirac work? Missing quantum geometry

## The “UFO” (fragmenting potential)

- Induces *surface Berry curvature* “in the sky,”  
*Not captured by the Dirac equation for the surface (Berry-flat)*
- Integrated surface Berry curvature:  $\theta(E) = \pi + \int_{0 \leq \varepsilon_{\mathbf{k}} \leq E} \Omega_{\mathbf{k}}$



# Why doesn't Dirac work? Missing quantum geometry

## The “UFO” (fragmenting potential)

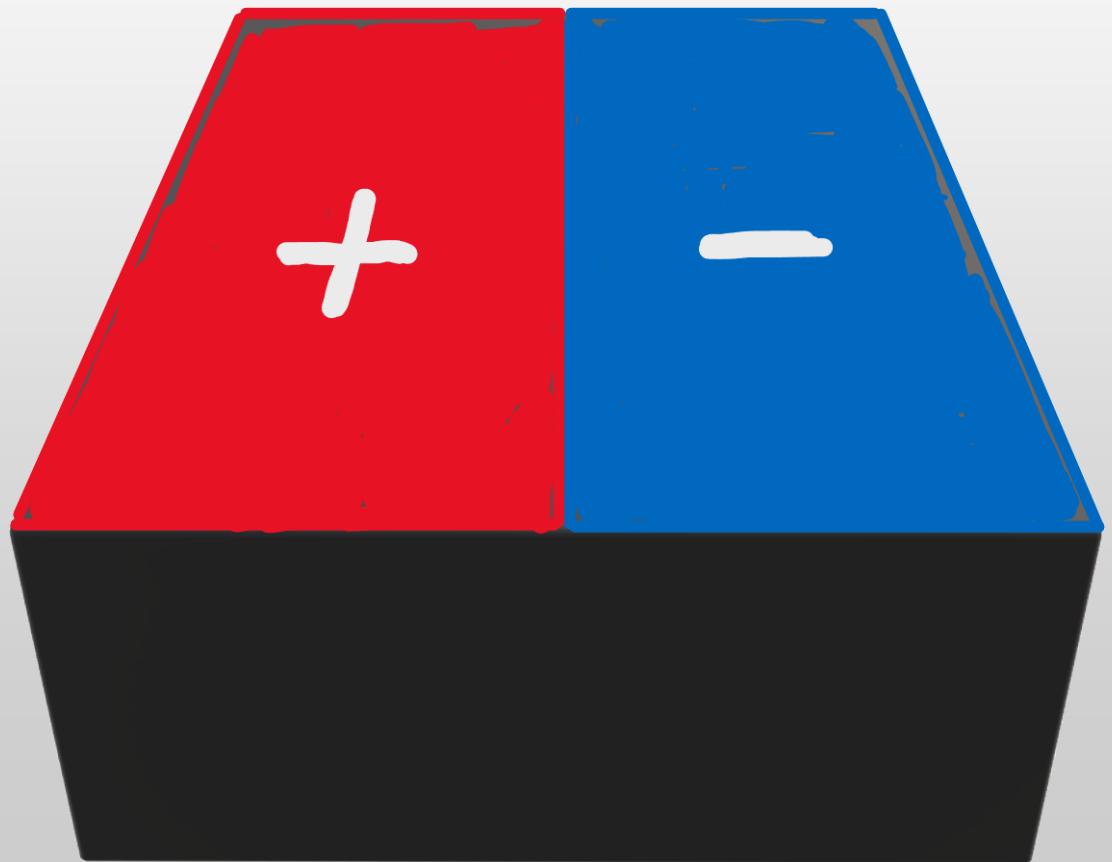
- Explains when and how spectrum-wide criticality can occur



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## The “UFO” (fragmenting potential)

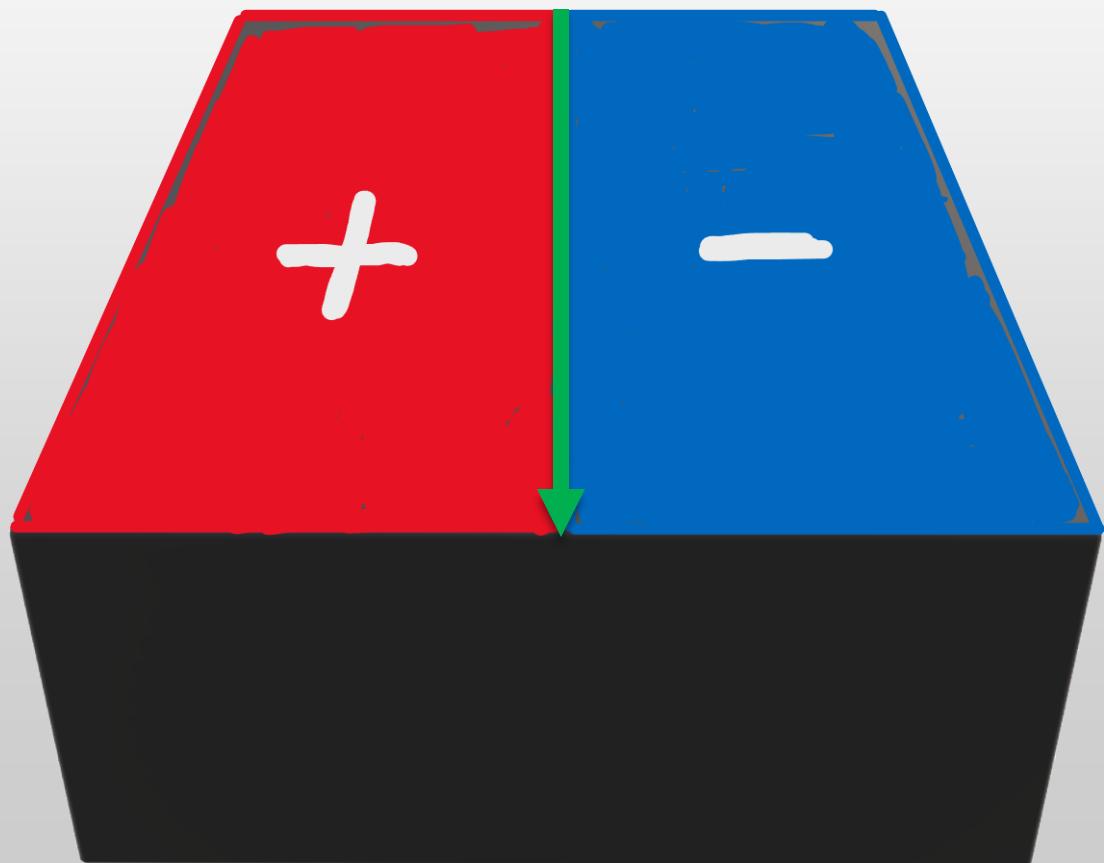
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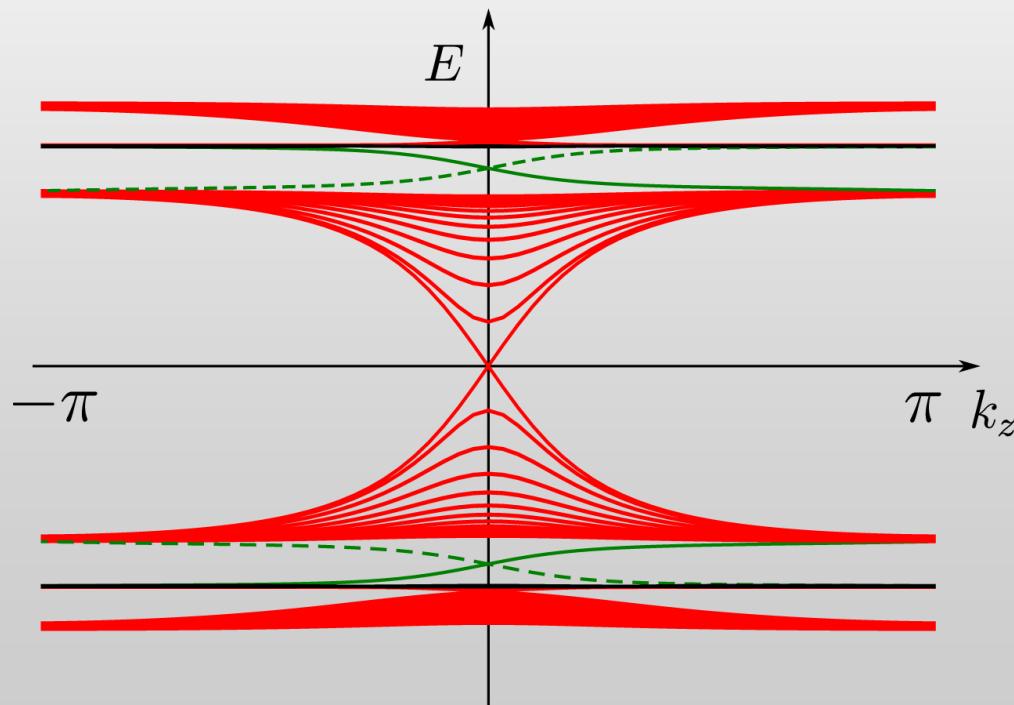
- Explains when and how spectrum-wide criticality can occur
- Chiral edge mode “in the sky!”



# Why doesn't Dirac work? Missing quantum geometry

## The “UFO” (fragmenting potential)

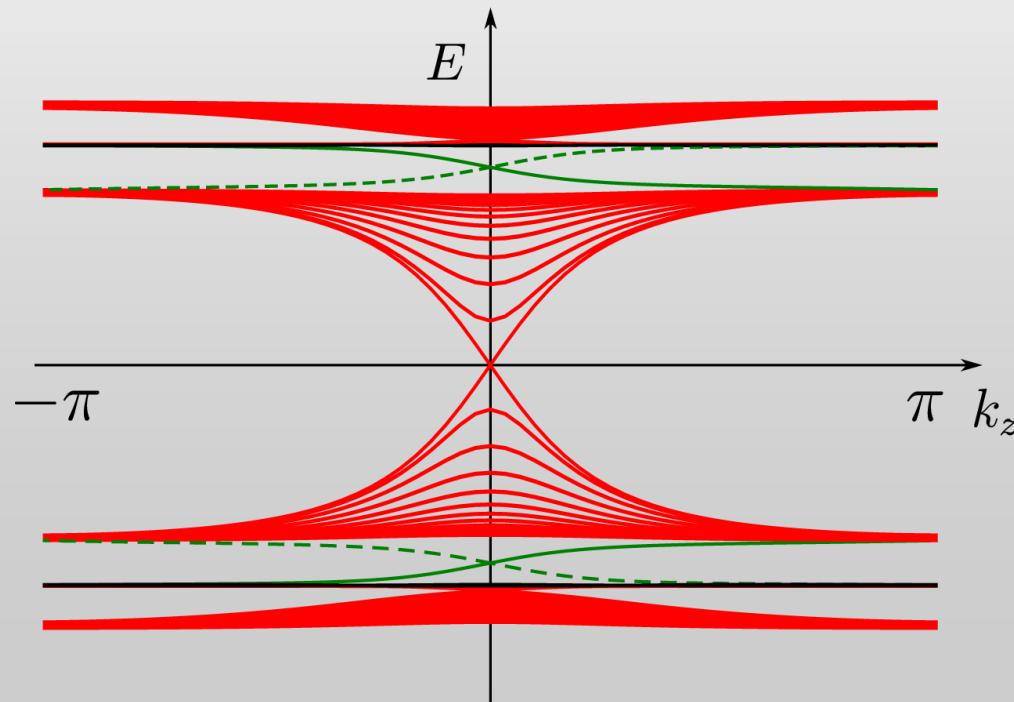
- Explains when and how spectrum-wide criticality can occur
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# Why doesn't Dirac work? Missing quantum geometry

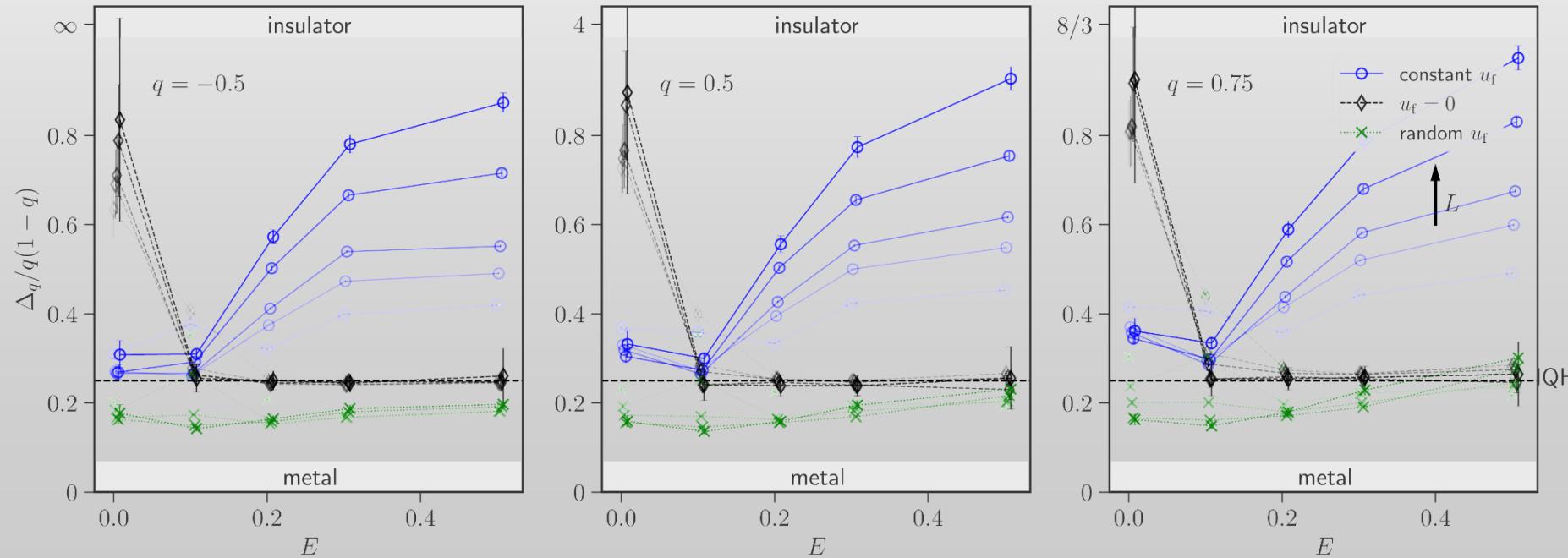
## The “UFO” (fragmenting potential)

- Explains when and how spectrum-wide criticality can occur
- Different signs of the potential introduce different surface domains of a “sky-Chern insulator”!
- 1D Chiral edge modes form at boundaries



# Why doesn't Dirac work? Missing quantum geometry

- Different signs of the potential introduce different surface domains of a “sky-Chern insulator”!
- 1D Chiral edge modes form at boundaries
- Spectrum-wide quantum criticality: When these edge modes percolate! (Random UFO with zero average...)

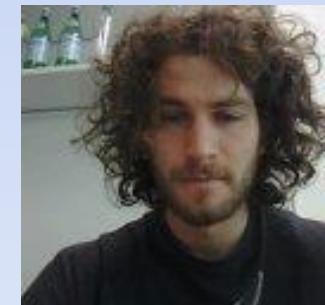
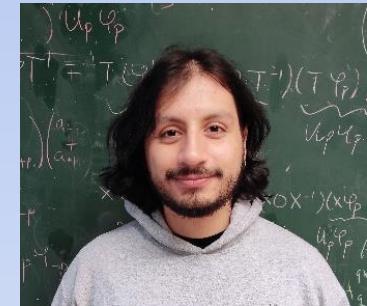
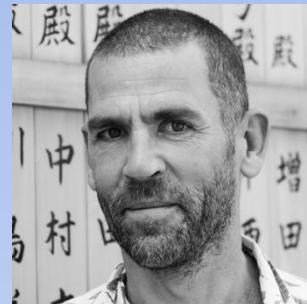


# Summary: Fragility of spectral flow in topological insulators

	class $d = 1$	$d = 2$	$d = 3$
A	0	$\mathbb{Z}^{\checkmark}$	0
AIII	$\mathbb{Z}^{\times}$	0	$\mathbb{Z}^{\times}$
AI	0	0	0
BDI	$\mathbb{Z}^{\times}$	0	0
D	$\mathbb{Z}_2^{\times}$	$\mathbb{Z}^{\checkmark}$	0
DIII	$\mathbb{Z}_2^{\times}$	$\mathbb{Z}_2^{\checkmark}$	$\mathbb{Z}^{\checkmark}/\times$
AII	0	$\mathbb{Z}_2^{\checkmark}$	$\mathbb{Z}_2^{\checkmark}$
CII	$2\mathbb{Z}^{\times}$	0	$\mathbb{Z}_2^{\times}$
C	0	$2\mathbb{Z}^{\checkmark}$	0
CI	0	0	$2\mathbb{Z}^{\times}$

$\checkmark$  : Not localizable

$\times$  : Localizable (e.g, UFO in 3D)



A. Altland, P. Brouwer, J. Dieplinger,  
M. S. Foster, M. Moreno-Gonzalez,  
and L. Trifunovic [arXiv:2308.12931](https://arxiv.org/abs/2308.12931)